

Deuxième école RMN du

GERM

Cargèse 2008

Analysis of micro-millisecond timescale motions of macromolecules in liquid state by Nuclear Magnetic Resonance

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Techniques to monitor protein dynamics

X-ray cristallography	B factors	time scales static disorder, crystal contacts,
X-Ray, neutron scattering Doniach, <i>Chem. Rev.</i> 2001, 101 ; Zacai, <i>science</i> 2000, 288.	size/shape modifications timescales (ps-ns) for ¹ H positions	
Fluorescence Weiss, Nat. Struct. Biol. 2000, 7; Yang, Science 2003, 302; Haustein, Curr. Opin. Struct. Biol. 2004, 14.	ensemble / single molecule cellular context	probes
Mass Spectroscopy (HX MS) radical footprinting Wales, Mass. Spectrom. Rev. 2006, 25 ; Busenlehner, Arch. Biochem. Biophys. 2005, 433. Guan, Trends. Biochem. Sci. 2005, 30.	large moldecular assemblies	
Mössbauer, Raman, 2D infrared spectroscopy		
Molecular dynamics		Forcefields Short timescales
NMR Boehr, <i>Chem. Rev.</i> 2006, 106 ,3055. Palmer, <i>Chem. Rev.</i> 2004, 104 , 3623.	 ⇒ 10⁻¹² ↔ 10⁵ s ⇒ Site-specific information ⇒ multiple atomic probes ¹H, ²H, ¹⁵N, ¹³C, ³¹P, ⇒ Simultaneous monitoring of probes ⇒ kinetic & termodynamic profile of dynamic processes 	 ⇒ isotope labeling ⇒ quantities ⇒ size limitation ⇒ complexity of the method ?

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Experimental aspects, applications

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How motions are visible in NMR?

Molecular motions influence NMR parameters Time scale of motions compared to NMR characteristic timescales

 \Rightarrow Three characteristic NMR timescales

Return of the spin system to equilibrium : T_1 / NMR signal lifetime : T_2

NMR experiment $\langle - \rangle$ spins system perturbation Liquid state: $T_1 \sim 100ms \leftrightarrow s$; $T_2 \sim 10ms \leftrightarrow s$

Minimal frequencies of motions that can be characterized during a single NMR experiment.

Spectral timescale : $\tau=1/\Delta v$

spectral features : chemical shifts range, couplings, ... $Hz \leftrightarrow kHz$

- Averaging of interactions by motions faster than the spectral dispersion due to these interactions.
- Perturbation of spectral features by motions in the same range than the spectral timescale.

Larmor timescale : ω_0

precession frequency of spins in a magnetic field ($B_0=\omega_0/\gamma$)

transitions efficiency between spin states is determined by molecular fluctuations (spectral densities fo motions) at these frequencies.

How motions are visible in NMR?

Molecular motions influence NMR parameters Time scale of motions compared to NMR characteristic timescales

Three characteristic NMR timescales
 Return of the spin system to equilibrium : T₁ / NMR signal lifetime : T₂
 Spectral timescale : τ=1/Δν
 Larmor timescale : ω₀





NMR timescales -1- Longitudinal relaxation time T_1

- √T₁ characterizes the time for a spin system to reach equilibrium.
- \checkmark Defines interscan delay
- ✓ Motions slower than T₁ cannot be characterize by a single spectra.
- √NB. T₁ depends on magnetic field
 B₀, nature of spins, molecule size,
 local flexibility, temperature, etc.

Dynamical processes slower than T_1

Real time folding Transient / out of equilibrium experiments Protein - ligand interactions H-N <-> D-N exchange rate constants



Wüthrich, « NMR of proteins and nucleic acids », Wiley Interscience, 1986

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```
holoP1 + apoP2 \iff apoP1 + holoP2
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magnetic field, molecular size and correlation time dependance of T_1 and T_2



 \checkmark Defined by the observed spectral width

✓ More precisely by the difference of resonance frequency between two spins from two different nuclei or from the same nuclei in two different states.

$$\Delta v (Hz) = \frac{\Delta \omega (ppm)}{2\pi} * 10^{-6} * \gamma B_0$$
$$\tau_{spect} = \frac{1}{\Delta v (Hz)}$$

 \checkmark Motions slower than $\Delta \nu$ have no effect on spectral feature

 $\checkmark \Delta v$ depends on magnetic field and on spins nature



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 $\checkmark \Delta v$ depends on the nature of the interactions between spins

Averaging of secular interactions by motions

 ✓ Dipolar interaction between spins magnetic moments.

$$\hat{H}_{IS}^{\text{DD}} = -\frac{\mu_0 \gamma_I \gamma_S \hbar}{4 \pi r_{IS}^3} \left(\frac{3 \cos^2 \theta_{IS} - 1}{2} \right) \left(3 \hat{I}_z \hat{S}_z - \hat{\mathbf{I}} \cdot \hat{\mathbf{S}} \right)$$
$$E/\hbar < 10^4 - 10^5 Hz$$

✓ Chemical shift anisotropy

$$\hat{H}_{I}^{CSA} = -\gamma_{S} \frac{c_{\parallel} - c_{\perp}}{3} B_{0} (3\cos^{2}\theta - 1) \hat{I}_{z}$$

$$E/\hbar < 10^{4} - 10^{5} Hz$$

- ✓ Interaction spin quadrupole/electric field $\hat{H}_{I}^{Q} \approx \omega_{I}^{Q} \left(3\hat{I}_{z}^{2} - \hat{\mathbf{I}}.\hat{\mathbf{I}}\right); \quad \omega_{I}^{Q} = \frac{3eQ_{I}}{4I(2I-1)}V_{zz}^{I}(\theta)$ $I > \frac{1}{2} \qquad E/\hbar \approx 2.10^{5}Hz(^{2}H) \text{ and } 3.10^{6}(^{14}N)$
- ✓ Unpaired electron

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motions generated to average interactions



NMR timescales -3- The "Larmor" timescale

 $\checkmark {\sf Defined}$ by the resonance frequencies of spins

√i.e. the energy difference between spins states levels

$$\omega_0 = -\gamma B_0$$

$$\tau_{Larmor} = \frac{1}{|\omega_0|} = \frac{1}{2\pi\nu_0}$$

- ✓ Motions in these timescale have no direct effect on spectra.
- ✓ Motions in these timescales are responsible for the efficiency of spins relaxation processes
- ✓ The relationship between motions and relaxation rate/time constants is not simple



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The method : relaxation dispersion NMR experiments

Quantitative evaluation of physical parameters of exchange from the relaxation dispersion data

The analysis depends of the timescale of the process

- is Fast exchange : $\Delta \omega \ll k_{ex}$
 - \Rightarrow One peak ω = $p_A \omega_A$ + $p_B \omega_B$
 - $\Rightarrow R_{ex} \propto B_0^2$
 - ⇒ Meiboom-Gill equation
 - \Rightarrow k_{ex}, Φ =p_Ap_B $\Delta \omega^{2}$, R_{inf}
- $\,$ Thtermediate exchange : $\Delta \omega$ ~ k_{ex}
 - \Rightarrow One or two broadened peaks
 - $\Rightarrow R_{ex} \propto B_0$
 - Carver-Richards equation
 - \Rightarrow k_{ex}, p_B, $\Delta \omega$, R_{inf}

 \square Slow exchange : $\Delta \square \gg k_{ex}$

- ⇒ Two peaks
- \Rightarrow R_{ex} independant of B₀
- Tollinger equation
- \Rightarrow ka, $\Delta \omega$, Rinf
- $\approx \mathsf{R}_{ex}^{(B)} \rightarrow \mathsf{k}_{BA} = p_A \mathsf{k}_{ex} ; \ \mathsf{R}_{ex}^{(A)} \rightarrow \mathsf{k}_{AB} = p_B \mathsf{k}_{ex} \\ \mathsf{R}_{ex}(B) \gg \ \mathsf{R}_{ex}(A)$
- Detection of minor populations from the broadenning of the major state resonance : use of Carver-Richards equation

 $p_A = 0.2$; $p_B = 0.8$ $R_2^{app} = R_2^{inf} + p_A p_B \Delta \omega^2 \tau_{ex} [1 - 4\tau_{ex} \nu_{CP} \tanh \frac{1}{4}$ $\frac{1}{4\tau_{ex}\nu_{CP}}$ $K_{ex}//N$ $R_{2}^{app} = R_{2}^{inf} + 0.5k_{ex} - \nu_{CP} \operatorname{acosh} (D_{+} \cosh \eta_{+} - D_{-} \cos \eta_{-})$ $D_{+} = 0.5 \left[1 + \frac{\phi + 2\Delta\omega^{2}}{\sqrt{\phi^{2} + \zeta^{2}}} \right] ; D_{-} = 0.5 \left| -1 + \frac{\phi + 2\Delta\omega^{2}}{\sqrt{\phi^{2} + \zeta^{2}}} \right|$ $\eta_{+} = \frac{\tau_{CP}}{\sqrt{(2)}} \sqrt{\phi + \sqrt{\phi^{2} + \zeta^{2}}}; \eta_{-} = \frac{\tau_{CP}}{\sqrt{(2)}} \sqrt{-\phi + \sqrt{\phi^{2} + \zeta^{2}}}$ $\phi = [(p_B - p_A)k_{er}]^2 - \Delta\omega^2 + 4p_A p_B k_{er}^2$ $\zeta = 2\Delta\omega(p_B - p_A)k_{er}$ $R_2^{app} = R_2^{inf} + k_a \left(1 - \frac{\sin \Delta \omega . \tau_{CP}}{\Delta \omega . \tau_{CP}} \right)$ -150 -50 50 -100 100 150 $\nu(Hz)$ $p_A v_A + p_B v_B$



The method : relaxation dispersion NMR experiments

Quantitative evaluation of physical parameters of exchange from the relaxation dispersion data









The method : relaxation dispersion NMR experiments relaxation dispersion experiments for the analysis of "invisible" states

An Isolated Helix Persists in a Sparsely Populated Form of KIX under Native Conditions



The method : relaxation dispersion NMR experiments

Quantitative evaluation of physical parameters of exchange from the relaxation dispersion data



pH=3.5

■ I34 ○ N42 ▼ L68 ◇ H71

(F)

R₂^{eff} (Hz)

200

400

600 800 v_{CPMG} (Hz) 1000

1200

1400



μ s-ms motions by NMR

conformational exchange processes and proteins function

⇒ Folding/unfolding processes

- ⇒ sequence/stability/flexibility relationship
- \Rightarrow low populated states

Transition state ensemble, minor folded/unfolded states

\Rightarrow Motions and function

- \Rightarrow enzymatic cycle
 - \Rightarrow low populated (high energy) states
- \Rightarrow regulation/signalization associated with conformational changes
- \Rightarrow interactions

μ s-ms motions and proteins function

- → Proteins : multiple conformers at thermal equilibrium in solution
 - > XRay cristallography : Selection of one conformation (low energy) in the cristallisation process
 - NMR : time averaged conformational ensemble
- → Ground state > 90%

→ High energy states, low populated : roles in substrates binding, enzymatic cycles, folding processes (transition states) ...



μ s-ms motions and proteins function : Practical course

Analysis of the relaxation dispersion pulse sequence

- The relaxation compensated CPMG scheme

The understanding of molecular mechanisms underlying proteins functions necessitates the characterization of the dynamics of the macromolecular systems, and of the minor transient states along the functional pathways.

Some recently published studies based on the use of relaxation dispersion experiments to investigate these questions will be discussed.